The discovery of the quantum Hall effect in 2D systems opens the door to topological phases of matter. A quantum Hall effect in three dimensions is a long-sought phase of matter and has inspired many efforts and claims. In this perspective, we review our proposal that guarantees a 3D quantum Hall effect. The proposal employs topologically protected Fermi arcs and ‘wormhole’ tunneling via the Weyl nodes in a 3D topological semimetal. The 1D edge states in this 3D quantum Hall effect show an example of $(d-2)$-dimensional boundary states. Possible signatures of the 3D quantum Hall effect have been observed in topological Dirac semimetals, but with many questions, which will attract more
research to verify the mechanism and realize the 3D quantum Hall effect in the future.

In a magnetic field, a moving charge feels a Lorentz force orthogonal to both its velocity and the magnetic field, leading to the Hall effect. Klaus von Klitzing discovered that in strong magnetic fields the Hall resistance of a 2D electron gas can be quantized into a series of plateaus in terms of $(e^2/h)n$ [1], where $e$ is the elementary charge, $h$ is Planck’s constant, and $n$ is an integer known as the ‘Chern number’ (named after the mathematician Shiing-Shen Chern). The quantum Hall effect has led to three Nobel Prizes in Physics (1985 von Klitzing; 1998 Tsui, Stormer, Laughlin; 2016 Thouless, Haldane, Kosterlitz). Usually, the quantum Hall effect takes place only in 2D systems. In a strong magnetic field, the energy spectrum of a 2D electron gas is quantized into Landau levels. The Landau levels deform at the sample edges and cross the Fermi energy, forming 1D edge states. Electrons can flow through the edge states without dissipation. When the Fermi energy is placed between two Landau levels, each edge state contributes a Hall conductance of $e^2/h$ and vanishing longitudinal conductance in the Hall-bar measurement. The quantization can be observed in two dimensions because the bulk states in the interior of the sample can be gapped. In contrast, a magnetic field quantizes the energy spectrum of a 3D electron gas into 1D Landau bands that disperse along the direction of the magnetic field. The dispersion prevents the quantization of the Hall conductance because the Fermi energy crosses some 1D Landau bands whose conductance is not quantized. Different schemes have been proposed to gap the 3D bulk states for the quantization of the Hall conductivity in three dimensions [2,3]. Nevertheless, a 3D quantum Hall effect remains a long-sought phase of matter [4–7].

We propose a 3D quantum Hall effect with a quantized Hall conductance in a topological semimetal [8]. The band structure of a topological semimetal looks like a 3D graphene [9–12], with the conduction and valence bands touching at the Weyl nodes (Fig. 1a). For momenta $k_z$ between the Weyl nodes, this band structure is equivalent to a 2D topological insulator, with topologically protected states on the surfaces (Fig. 1b–d) parallel to the $z$ direction. The Fermi surface of the surface states is known as the Fermi arcs (red and blue curves in Fig. 1a–d). The Fermi-arc surface states form a unique 2D electron gas, half from the top surface and half from the bottom surface (Fig. 1c and d). It may host a quantum Hall effect. If there were only the top surface (Fig. 1g), the Fermi-arc surface states could not support a complete cyclotron motion in real space (Fig. 1f); then there would be no Landau levels, edge states, or quantum Hall effect. Fortunately, the top and bottom surfaces can form a complete 2D electron gas, with a closed Fermi surface connected by the Weyl nodes. Driven by the $y$-direction magnetic field, an electron performs half of a cyclotron motion on the top Fermi arc, then tunnels via a Weyl node to the bottom Fermi arc to complete the cyclotron motion. In this way, the top and bottom Fermi arcs together support a complete cyclotron motion and the quantum Hall effect. More importantly, the Weyl nodes are 3D singularities in momentum space, so according to the uncertainty principle they can connect 2D surfaces separated infinitely far apart in real space. This is why we call...
it a 3D quantum Hall effect. This is like the wormhole effect, which connects 3D spaces via higher-dimensional singularities.

Recently, quantized Hall resistance plateaus have been experimentally observed in the topological semimetal Cd$_3$As$_2$ [13–15], with thickness ranging from 10 to 80 nm. They cannot be regarded as 2D. Nevertheless, several questions still hold. First, Cd$_3$As$_2$ is a Dirac semimetal, composed of two time-reversed Weyl semimetals. At a single surface, there is a complete 2D electron gas, formed by two time-reversed half-2D electron gases of the Fermi-arc surface states. There may also be a trivial quantum Hall effect on a single surface. Second, the 3D bulk states quantize 2D subbands for those thicknesses. If the 3D bulk states cannot be depleted entirely, they also have a trivial quantum Hall effect. These two issues may explain the 2-fold and 4-fold degenerate Hall resistance plateaus observed in the experiments. To deplete the 3D bulk states, the Fermi energy has to be placed exactly at the Weyl nodes. Clarifying how to distinguish these trivial mechanisms from the 3D quantum Hall effect will be an interesting direction. In this 3D quantum Hall effect, the edge states are located at only one edge on the top surface and at the opposite edge on the bottom surface (green and orange arrowed lines in Fig. 1d and e), which can be probed by scanning tunneling microscopy. The 3D quantum Hall effect may be realized in other systems with novel surface states. More research will be necessary to verify the mechanism and realize the 3D quantum Hall effect in the future.

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